## Variational Methods \& Optimal Control

lecture 12

Matthew Roughan
[matthew.roughan@adelaide.edu.au](mailto:matthew.roughan@adelaide.edu.au)
Discipline of Applied Mathematics
School of Mathematical Sciences
University of Adelaide

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## Numerical Solutions

The E-L equations may be hard to solve
Natural response is to find numerical methods

- Numerical solution of E-L DE
$\triangleright$ we won't consider these here (see other courses)
- Euler's finite difference method
- Ritz (Rayleigh-Ritz)
$\triangleright$ In 2D: Kantorovich's method


## Euler's finite difference method

We can approximate our function (and hence the integral) onto a finite grid. In this case, the problem reduces to a standard multivariable maximization (or minimization) problem, and we find the solution by setting the derivatives to zero. In the limit as the grid gets finer, this approximates the E-L equations.

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## Numerical Approximation

Numerical approximation of integrals:

- use an arbitrary set of mesh points $a=x_{0}<x_{1}<x_{2}<\cdots<x_{n}=b$.
- approximate

$$
y^{\prime}\left(x_{i}\right)=\frac{y_{i+1}-y_{i}}{x_{i+1}-x_{i}}=\frac{\Delta y_{i}}{\Delta x_{i}}
$$

- rectangle rule

$$
F\{y\}=\int_{a}^{b} f\left(x, y, y^{\prime}\right) d x \simeq \sum_{i=0}^{n-1} f\left(x_{i}, y_{i}, \frac{\Delta y_{i}}{\Delta x_{i}}\right) \Delta x_{i}=\bar{F}(\mathbf{y})
$$

$\bar{F}(\cdot)$ is a function of the vector $\mathbf{y}=\left(y_{1}, y_{2}, \ldots, y_{n}\right)$.

## Finite Difference Method (FDM)

Treat this as a maximization of a function of $n$ variables, so that we require

$$
\frac{\partial \bar{F}}{\partial y_{i}}=0
$$

for all $i=1,2, \ldots, n$.
Typically use uniform grid so $\Delta x_{i}=\Delta x=(b-a) / n$.

## Simple Example

Find extremals for

$$
F\{y\}=\int_{0}^{1}\left[\frac{1}{2} y^{\prime 2}+\frac{1}{2} y^{2}-y\right] d x
$$

with $y(0)=0$ and $y(1)=0$.
E-L equations $y^{\prime \prime}-y=1$.

## Simple Example: direct solution

E-L equations $y^{\prime \prime}-y=-1$
Solution to homogeneous equations $y^{\prime \prime}-y=0$ is given by $e^{\lambda x}$ giving
characteristic equation $\lambda^{2}-1=0$, so $\lambda= \pm 1$.
Particular solution $y=1$
Final solution is

$$
y(x)=A e^{x}+B e^{-x}+1
$$

The boundary conditions $y(0)=y(1)=0$ constrain $A+B=-1$ and $A e+B e^{-1}=-1$, so $A e+(1-A) e^{-1}=1$, so $A=\frac{e^{-1}-1}{e-e^{-1}}$ and $B=\frac{1-e}{e-e^{-1}}$.
Then the exact solution to the extremal problem is

$$
y(x)=\frac{e^{-1}-1}{e-e^{-1}} e^{x}+\frac{1-e}{e-e^{-1}} e^{-x}-1
$$

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## Simple Example: Euler's FDM

Find extremals for

$$
F\{y\}=\int_{0}^{1}\left[\frac{1}{2} y^{\prime 2}+\frac{1}{2} y^{2}-y\right] d x
$$

Euler's FDM.

- Take the grid $x_{i}=i / n$, for $i=0,1, \ldots, n$ so
$\triangleright$ end points $y_{0}=0$ and $y_{n}=0$
$\triangleright \Delta x=1 / n$
$\triangleright \Delta y_{i}=y_{i+1}-y_{i}$
- So
$\triangleright y_{i}^{\prime}=\Delta y_{i} / \Delta x=n\left(y_{i+1}-y_{i}\right)$
$\triangleright$ and

$$
y_{i}^{\prime 2}=n^{2}\left(y_{i}^{2}-2 y_{i} y_{i+1}+y_{i+1}^{2}\right)
$$

## Simple Example: Euler's FDM

Find extremals for

$$
F\{y\}=\int_{0}^{1}\left[\frac{1}{2} y^{\prime 2}+\frac{1}{2} y^{2}-y\right] d x
$$

Its FDM approximation is

$$
\begin{aligned}
\bar{F}(\mathbf{y}) & =\sum_{i=0}^{n-1} f\left(x_{i}, y_{i}, y_{i}^{\prime}\right) \Delta x \\
& =\sum_{i=0}^{n-1} \frac{1}{2} n^{2}\left(y_{i}^{2}-2 y_{i} y_{i+1}+y_{i+1}^{2}\right) \Delta x+\left(y_{i}^{2} / 2-y_{i}\right) \Delta x \\
& =\sum_{i=0}^{n-1} \frac{1}{2} n\left(y_{i}^{2}-2 y_{i} y_{i+1}+y_{i+1}^{2}\right)+\frac{y_{i}^{2} / 2-y_{i}}{n}
\end{aligned}
$$

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## Simple Example: end-conditions

- We know the end conditions $y(0)=y(1)=0$, which imply that

$$
y_{0}=y_{n}=0
$$

- Include them into the objective using Lagrange multipliers

$$
\bar{H}(\mathbf{y})=\sum_{i=0}^{n-1} \frac{1}{2} n\left(y_{i}^{2}-2 y_{i} y_{i+1}+y_{i+1}^{2}\right)+\frac{y_{i}^{2} / 2-y_{i}}{n}+\lambda_{0} y_{0}+\lambda_{n} y_{n}
$$

## Simple Example: Euler's FDM

Taking derivatives, note that $y_{i}$ only appears in two terms of the FDM approximation

$$
\begin{aligned}
\bar{H}(\mathbf{y}) & =\sum_{i=0}^{n-1} \frac{1}{2} n\left(y_{i}^{2}-2 y_{i} y_{i+1}+y_{i+1}^{2}\right)+\frac{y_{i}^{2} / 2-y_{i}}{n}+\lambda_{0} y_{0}+\lambda_{n} y_{n} \\
\frac{\partial \bar{H}(\mathbf{y})}{\partial y_{i}} & = \begin{cases}n\left(y_{0}-y_{1}\right)+\frac{y_{0}-1}{n}+\lambda_{0} & \text { for } i=0 \\
n\left(2 y_{i}-y_{i+1}-y_{i-1}\right)+\frac{y_{i}}{n}-\frac{1}{n} & \text { for } i=1, \ldots, n-1 \\
n\left(y_{n}-y_{n-1}\right)+\lambda_{n} & \text { for } i=n\end{cases}
\end{aligned}
$$

We need to set the derivatives to all be zero, so we now have $n+3$ linear equations, including $y_{0}=y_{n}=0$, and $n+3$ variables including the two Lagrange multipliers. We can solve this system numerically using, e.g., matlab.

$$
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$$

Simple Example: Euler's FDM
Example: $n=4$, solve

$$
A \mathbf{z}=\mathbf{b}
$$

where
$A=\left(\begin{array}{rrrrr}-4.00 & & & & -4.00 \\ 8.25 & -4.00 & & & \\ -4.00 & 8.25 & -4.00 & & \\ & -4.00 & 8.25 & -4.00 & \\ & & -4.00 & 8.25 & -4.00 \\ & & -4.00 & 8.25 & -4.00\end{array}\right)$ and $\mathbf{b}=\left(\begin{array}{l}0.00 \\ 0.25 \\ 0.25 \\ 0.25 \\ 0.25 \\ 0.00 \\ 0.00\end{array}\right)$

- first $n+1$ terms of $\mathbf{z}$ give $\mathbf{y}$
- last two terms given the Lagrange multipliers $\lambda_{0}$ and $\lambda_{n}$


## Simple example: results

## Convergence of Euler's FDM

The condition for a stationary point becomes

$$
\frac{\partial \bar{F}}{\partial y_{i}}=\frac{\partial f}{\partial y_{i}}\left(x_{i}, y_{i}, y_{i}^{\prime}\right)-\frac{\frac{\partial f}{\partial y_{i}^{\prime}}\left(x_{i}, y_{i}, \frac{\Delta y_{i}}{\Delta x}\right)-\frac{\partial f}{\partial y_{i}^{\prime}}\left(x_{i-1}, y_{i-1}, \frac{\Delta y_{i-1}}{\Delta x}\right)}{\Delta x}=0
$$

In limit $n \rightarrow \infty$, then $\Delta x \rightarrow 0$, and so we get

$$
\frac{\partial f}{\partial y}-\frac{d}{d x}\left(\frac{\partial f}{\partial y^{\prime}}\right)=0
$$

which are the Euler-Lagrange equations.

- i.e., the finite difference solution converges to the solution of the E-L equations

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## Convergence of Euler's FDM

$$
\bar{F}(\mathbf{y})=\sum_{i=0}^{n-1} f\left(x_{i}, y_{i}, \frac{\Delta y_{i}}{\Delta x}\right) \Delta x \quad \text { and } \quad \Delta y_{i}=y_{i+1}-y_{i}
$$

Only and two terms in the sum involve $y_{i}$, so

$$
\begin{aligned}
\frac{\partial \bar{F}}{\partial y_{i}}= & \frac{\partial}{\partial y_{i}} f\left(x_{i-1}, y_{i-1}, \frac{\Delta y_{i-1}}{\Delta x}\right)+\frac{\partial}{\partial y_{i}} f\left(x_{i}, y_{i}, \frac{\Delta y_{i}}{\Delta x}\right) \\
= & \frac{1}{\Delta x} \frac{\partial f}{\partial y_{i}^{\prime}}\left(x_{i-1}, y_{i-1}, \frac{\Delta y_{i-1}}{\Delta x}\right) \\
& +\frac{\partial f}{\partial y_{i}}\left(x_{i}, y_{i}, \frac{\Delta y_{i}}{\Delta x}\right)-\frac{1}{\Delta x} \frac{\partial f}{\partial y_{i}^{\prime}}\left(x_{i}, y_{i}, \frac{\Delta y_{i}}{\Delta x}\right) \\
= & \frac{\partial f}{\partial y_{i}}\left(x_{i}, y_{i}, y_{i}^{\prime}\right)-\frac{\frac{\partial f}{\partial y_{i}^{\prime}}\left(x_{i}, y_{i}, \frac{\Delta y_{i}}{\Delta x}\right)-\frac{\partial f}{\partial y_{i}^{\prime}}\left(x_{i-1}, y_{i-1}, \frac{\Delta y_{i-1}}{\Delta x}\right)}{\Delta x}
\end{aligned}
$$

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## Comments

- There are lots of ways to improve Euler's FDM
$\triangleright$ use a better method of numerical quadrature (integration)
$\star$ trapezoidal rule
* Simpson's rule
$\star$ Romberg's method
$\triangleright$ use a non-uniform grid
$\star$ make it finer where there is more variation
- We can use a different approach that can be even better


## Ritz's method

In Ritz's method (called Kantorovich's methods where there is more than one independent variable), we approximate our functions (the extremal in particular) using a family of simple functions. Again we can reduce the problem into a standard multivariable maximization problem, but now we seek coefficients for our approximation.

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## Ritz's method

Assume we can approximate $y(x)$ by

$$
y(x)=\phi_{0}(x)+c_{1} \phi_{1}(x)+c_{2} \phi_{2}(x)+\cdots+c_{n} \phi_{n}(x)
$$

where we choose a convenient set of functions $\phi_{j}(x)$ and find the values of $c_{j}$ which produce an extremal.

For fixed end-point problem:

- Choose $\phi_{0}(x)$ to satisfy the end conditions.
- Then $\phi_{j}\left(x_{0}\right)=\phi_{j}\left(x_{1}\right)=0$ for $j=1,2, \ldots, n$

The $\phi$ can be chosen from standard sets of functions, e.g. power series, trigonometric functions, Bessel functions, etc. (but must be linearly independent)

## Ritz's method

- select $\left\{\phi_{j}\right\}_{j=0}^{n}$
- Approximate $y_{n}(x)=\phi_{0}(x)+c_{1} \phi_{1}(x)+c_{2} \phi_{2}(x)+\cdots+c_{n} \phi_{n}(x)$
- Approximate $F\{y\} \simeq F\left\{y_{n}\right\}=\int_{x_{0}}^{x_{1}} f\left(x, y_{n}, y_{n}^{\prime}\right) d x$
- Integrate to get $F\left\{y_{n}\right\}=F_{n}\left(c_{1}, c_{2}, \ldots, c_{n}\right)$
- $F_{n}$ is a known function of $n$ variables, so we can maximize (or minimize) it as usual by

$$
\frac{\partial F_{n}}{\partial c_{i}}=0
$$

for all $i=1,2, \ldots, n$.

## Upper bounds

Assume the extremal of interest is a minimum, then for the extremal

$$
F\{y\}<F\{\hat{y}\}
$$

for all $\hat{y}$ within the neighborhood of $y$. Assume our approximating function $y_{n}$ is close enough to be in that neighborhood, then

$$
F\{y\} \leq F\left\{y_{n}\right\}=F_{n}(\mathbf{c})
$$

so the approximation provides an upper bound on the minimum $F\{y\}$. Another way to think about it is that we optimize on a smaller set of possible functions $y$, so we can't get quite as good a minimum.

## Simple Example

Find extremals for

$$
F\{y\}=\int_{0}^{1}\left[\frac{1}{2} y^{\prime 2}+\frac{1}{2} y^{2}-y\right] d x
$$

with $y(0)=0$ and $y(1)=0$.
E-L equations $y^{\prime \prime}-y=1$, but we shall bypass the E-L equations to use Ritz's method.

$$
y_{n}(x)=\phi_{0}(x)+\sum_{i=1}^{n} c_{i} \phi_{i}(x)
$$

where we take $\phi_{0}(x)=0$ and $\phi_{i}(x)=x^{i}(1-x)^{i}$.

## Simple Example

Simple approximation $y_{1}=c_{1} \phi_{1}(x)$ we get

$$
F_{1}\left(c_{1}\right)=F\left\{y_{1}\right\}=\int_{0}^{1}\left[\frac{1}{2} c_{1}^{2} \phi_{1}^{\prime 2}+c_{1}^{2} \frac{1}{2} \phi_{1}^{2}-c_{1} \phi_{1}\right] d x
$$

Now $\phi(x)=x(1-x)$ so $\phi_{1}^{\prime}=1-2 x$, and

$$
\begin{aligned}
F_{1}\left(c_{1}\right) & =\int_{0}^{1}\left[\frac{c_{1}^{2}}{2}(1-2 x)^{2}+\frac{c_{1}^{2}}{2} x^{2}(1-x)^{2}-c_{1} x(1-x)\right] d x \\
& =\frac{c_{1}^{2}}{2} \int_{0}^{1}\left[1-4 x+5 x^{2}-x^{4}\right] d x+c_{1} \int_{0}^{1}\left[-x+x^{2}\right] d x \\
& =\frac{c_{1}^{2}}{2}\left[x-2 x^{2}+5 x^{3} / 3-x^{5} / 5\right]_{0}^{1}+c_{1}\left[-x^{2} / 2+x^{3} / 3\right]_{0}^{1} \\
& =\frac{c_{1}^{2}}{2} \frac{11}{30}-\frac{c_{1}}{6}
\end{aligned}
$$

## Simple Example

We solve for $c_{1}$ by setting

$$
\frac{d F_{1}}{d c_{1}}=\frac{11 c_{1}}{30}-\frac{1}{6}=0
$$

to get $c_{1}=5 / 11$, so the approximate extremal is

$$
y_{1}(x)=\frac{5}{11} x(1-x)
$$

The value of the approximate functional at this point is

$$
F_{1}(5 / 11)=\frac{c_{1}^{2}}{2} \frac{11}{30}-\frac{c_{1}}{6}=-0.37879
$$

which is an upper bound on the true value of the functional on the extremal.

Simple example: results


## Alternate approach

Choose $\phi_{1}(x)=\sin (\pi x)$ (use the first element of a trigonometric series to approximate $y$ ). Then, $\phi^{\prime}(x)=\pi \cos (\pi x)$, and so the functional is

$$
\begin{aligned}
F_{1}\left(c_{1}\right) & =F\left\{c_{1} \phi_{1}\right\}=\int_{0}^{1}\left[\frac{1}{2} c_{1}^{2} \phi_{1}^{\prime 2}+c_{1}^{2} \frac{1}{2} \phi_{1}^{2}-c_{1} \phi_{1}\right] d x \\
& =\int_{0}^{1}\left[\frac{c_{1}^{2} \pi^{2}}{2} \cos ^{2}(\pi x)+\frac{c_{1}^{2}}{2} \sin ^{2}(\pi x)-c_{1} \sin (\pi x)\right] d x
\end{aligned}
$$

Now $\int_{0}^{1} \cos ^{2}(\pi x)=\int_{0}^{1} \sin ^{2}(\pi x)=1 / 2$,
and $\int_{0}^{1} \sin (\pi x)=\left[-\frac{1}{\pi} \cos (\pi x)\right]_{0}^{1}=-2 / \pi$, so

$$
F\left(c_{1}\right)=\frac{c_{1}^{2}}{2} \frac{1}{2}\left[\pi^{2}+1\right]-\frac{2}{\pi} c_{1}
$$

## Alternate approach

Once again we solve for $c_{1}$ by setting

$$
\frac{d F_{1}}{d c_{1}}=c_{1} \frac{1}{2}\left[\pi^{2}+1\right]-\frac{2}{\pi}=0
$$

to get $c_{1}=\frac{4}{\pi\left(\pi^{2}+1\right)}$, so the approximate extremal is

$$
y_{1}(x)=\frac{4}{\pi\left(\pi^{2}+1\right)} \sin (\pi x)
$$

