# Variational Methods \& Optimal Control 

lecture 25

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## Conservation Laws

One of the more exciting things we can derive relates to fundamental physics laws: conservation of energy, momentum, and angular momentum. We can now derive all of these from an underlying principle: Noether's theorem.

## Hamilton's principle

We now have a group of equivalent methods

- Euler-Lagrange equations
- Hamilton's equations
- Hamilton-Jacobi equation

We saw earlier that these can give us other methods

- Hamilton's principle $\Rightarrow$ Newton's laws of motion
- When $L$ is not explicitly dependent on $t$, then the Hamiltonian $H$ is constant in time.
$\triangleright$ conservation of energy
$\triangleright$ this is an illustration of a symmetry in the problem appearing in the Hamiltonian

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## Conservation laws

Given the functional

$$
F\{y\}=\int_{x_{0}}^{x_{1}} f\left(x, y, y^{\prime}, \ldots, y^{(n)}\right) d x
$$

if there is a function $\phi\left(x, y, y^{\prime}, \ldots, y^{(k)}\right)$ such that

$$
\frac{d}{d x} \phi\left(x, y, y^{\prime}, \ldots, y^{(k)}\right)=0
$$

for all extremals of $F$, then this is called a $k$ th order conservation law

- use obvious extension for functionals of several dependent variables.


## Conservation law example

Given the functional

$$
F\{y\}=\int_{x_{0}}^{x_{1}} f\left(y, y^{\prime}\right) d x
$$

where $f$ is not explicitly dependent on $t$, we know that the Hamiltonian

$$
H=y^{\prime} \frac{\partial f}{\partial y^{\prime}}-f
$$

is constant, and so

$$
\frac{d H}{d x}=0
$$

is a first order conservation law for the system.

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## Several independent variables

For functionals of several independent variables, e.g.

$$
F\{z\}=\iint_{\Omega} z(x, y) d x d y
$$

the equivalent conservation law is

$$
\nabla \cdot \phi=0
$$

For some function $\phi\left(x, y, z, z^{\prime}, \ldots, z^{(k)}\right)$.

- Results here can be extended to these cases, but we won't look at them here.


## Conservation laws

- physically interesting
$\triangleright$ tell you about system of interest
- can simplify solution
$\triangleright \phi\left(x, y, y^{\prime}, \ldots, y^{(k)}\right)=$ const is an order $k$ DE, rather than E-L equations which are order $2 n$
- $\phi\left(x, y, y^{\prime}, \ldots, y^{(k)}\right)=$ const is often called the first integral of the

E-L equations
$\triangleright$ RHS is a constant of integration (determined by boundary conditions)

- how do we find conservation laws?
$\triangleright$ Noether's theorem

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## Variational symmetries

The key to finding conservation laws lies in finding symmetries in the problem.

- "symmetries" are the result of transformations under which the functional is invariant
- E.G. time invariance symmetry results in constant $H$
- more generally, take a parameterized family of smooth transforms

$$
X=\theta(x, y ; \varepsilon), \quad Y=\phi(x, y ; \varepsilon)
$$

where

$$
x=\theta(x, y ; 0), \quad y=\phi(x, y ; 0)
$$

e.g. we get the identity transform for $\varepsilon=0$

- examples translations and rotations


## Jacobian

The Jacobian is

$$
J=\left|\begin{array}{cc}
\theta_{x} & \theta_{y} \\
\phi_{x} & \phi_{y}
\end{array}\right|=\theta_{x} \phi_{y}-\theta_{y} \phi_{x}
$$

- smooth: if functions $x$ and $y$ have continuous partial derivatives.
- non-singular: if Jacobian is non-zero (and hence an inverse transform exists)

Now for $\varepsilon=0$, we require the identity transform, so $J=1$. Also, we require a smooth transform, so $J$ is a smooth function of $\varepsilon$, and so for sufficiently small $|\varepsilon|$, the transform is non-singular.

## Example transformations

- translations ( $\varepsilon$ is the translation distance)


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## Example transformations

- translations ( $\varepsilon$ is the translation distance)



## Example transformations

- rotations ( $\varepsilon$ is the rotation angle)

$$
X=x \cos \varepsilon+y \sin \varepsilon \quad Y=-x \sin \varepsilon+y \cos \varepsilon
$$

has Jacobian

$$
J=\cos ^{2} \varepsilon+\sin ^{2} \varepsilon=1
$$

and inverse

$$
x=X \cos \varepsilon-Y \sin \varepsilon \quad y=X \sin \varepsilon+Y \cos \varepsilon
$$

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## Example transformations

- rotations ( $\varepsilon$ is the rotation angle)



## Example transformations

- rotations ( $\varepsilon$ is the rotation angle)

$$
X=x \cos \varepsilon+y \sin \varepsilon \quad Y=-x \sin \varepsilon+y \cos \varepsilon
$$

To derive this, change coordinates to polar coordinates

$$
x=r \cos (\theta) \quad \text { and } \quad y=r \sin (\theta)
$$

Under a rotation by $\varepsilon$, the new coordinates $(X, Y)$ are

$$
X=r \cos (\theta-\varepsilon) \quad \text { and } \quad Y=r \sin (\theta-\varepsilon)
$$

Use trig. identities $\cos (u-v)=\cos u \cos v+\sin u \sin v$ and $\sin (u-v)=\sin u \cos v-\cos u \sin v$, to get

$$
X=r \cos (\theta) \cos (\varepsilon)+r \sin (\theta) \sin (\varepsilon)=x \cos (\varepsilon)+y \sin (\varepsilon)
$$

$$
Y=r \sin (\theta) \cos (\varepsilon)-r \cos (\theta) \sin (\varepsilon)=y \cos (\varepsilon)-x \sin (\varepsilon)
$$

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## Transformation of a function

Given a function $y(x)$, we can rewrite $Y(X)$ using the inverse transformation, e.g.

$$
\phi^{-1}(X, Y(X) ; \varepsilon)=y(x)=y\left(\theta^{-1}(X, Y ; \varepsilon)\right)
$$

For example, taking the curve $y=x$ under rotations

$$
X \sin \varepsilon+Y \cos \varepsilon=X \cos \varepsilon-Y \sin \varepsilon
$$

which we rearrange to get

$$
Y(X)=\frac{\cos \varepsilon-\sin \varepsilon}{\cos \varepsilon+\sin \varepsilon} X
$$

Similarly we can derive $Y^{\prime}(X)$

## Transform invariance

## If

$$
\int_{x_{0}}^{x_{1}} f\left(x, y, y^{\prime}(x)\right) d x=\int_{X_{0}}^{X_{1}} f\left(X, Y, Y^{\prime}(X)\right) d X
$$

for all smooth functions $y(x)$ on $\left[x_{0}, x_{1}\right]$ then we say that the functional in invariant under the transformation.

- also called variational invariance
- The transform is called a variational symmetry
- Related to conservation laws

Also note that the E-L equations are invariant under such a transform, e.g. they produce the same extremal curves.

## Examples

- translations:

$$
\begin{aligned}
(X, Y) & =(x+\varepsilon, y) \Rightarrow(\xi, \eta)=(1,0) \\
\text { or } \quad \Rightarrow \quad(X, Y) & =(x, y+\varepsilon) \quad \Rightarrow \quad(\xi, \eta)=(0,1)
\end{aligned}
$$

- rotations:
$X=\theta(x, y ; \varepsilon)=x \cos \varepsilon+y \sin \varepsilon \quad Y=\phi(x, y ; \varepsilon)=-x \sin \varepsilon+y \cos \varepsilon$
So

$$
\begin{aligned}
\xi & =\left.\frac{\partial \theta}{\partial \varepsilon}\right|_{\varepsilon=0}=-x \sin \varepsilon+\left.y \cos \varepsilon\right|_{\varepsilon=0}=y \\
\eta & =\left.\frac{\partial \phi}{\partial \varepsilon}\right|_{\varepsilon=0}=-x \cos \varepsilon-\left.y \sin \varepsilon\right|_{\varepsilon=0}=-x
\end{aligned}
$$

## Infinitesimal generators

For small $\varepsilon$ we can use Taylor's theorem to write

$$
\begin{aligned}
X & =\theta(x, y ; 0)+\left.\varepsilon \frac{\partial \theta}{\partial \varepsilon}\right|_{(x, y ; 0)}+O\left(\varepsilon^{2}\right) \\
Y & =\phi(x, y ; 0)+\left.\varepsilon \frac{\partial \phi}{\partial \varepsilon}\right|_{(x, y ; 0)}+O\left(\varepsilon^{2}\right)
\end{aligned}
$$

Define the infinitesimal generators

$$
\xi(x, y)=\left.\frac{\partial \theta}{\partial \varepsilon}\right|_{(x, y ; 0)} \quad \eta(x, y)=\left.\frac{\partial \phi}{\partial \varepsilon}\right|_{(x, y ; 0)}
$$

and then for small $\varepsilon$

$$
\begin{aligned}
& X \simeq x+\varepsilon \xi \\
& Y \simeq y+\varepsilon \eta
\end{aligned}
$$

- Amalie Emmy Noether, 23 March 1882 - 14 April 1935
- Described by Einstein and many others as the most important woman in the history of mathematics.
- Most of her work was in algebra
- Worked at the Mathematical Institute of Erlangen without pay for seven years
- Invited by David Hilbert and Felix Klein to join the mathematics department at the University of Göttingen, a world-renowned center of mathematical research. The philosophicqal faculty objected, however, and she spent four years lecturing under Hilbert's name.


## Noether's theorem

Suppose the $f\left(x, y, y^{\prime}\right)$ is variationally invariant on $\left[x_{0}, x_{1}\right]$ under a transform with infinitesimal generators $\xi$ and $\eta$, then

$$
\eta p-\xi H=\mathrm{const}
$$

along any extremal of

$$
F\{y\}=\int_{x_{0}}^{x_{1}} f\left(x, y, y^{\prime}\right) d x
$$

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## Example (i)

Invariance in translations in $x$, i.e.

$$
\begin{aligned}
(X, Y) & =(x+\varepsilon, y) \\
(\xi, \eta) & =(1,0)
\end{aligned}
$$

So, a system with such invariance has

$$
H=\text { const }
$$

which is what we showed earlier regarding functionals with no explicit dependence on $x$.

## Example (ii)

Invariance in translations in $y$, i.e.

$$
\begin{aligned}
(X, Y) & =(x, y+\varepsilon) \\
(\xi, \eta) & =(0,1)
\end{aligned}
$$

So, a system with such invariance has

$$
p=\text { const }
$$

which is what we showed earlier regarding functionals with no explicit dependence on $y$.

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## More than one dependent variable

Transforms with more than one dependent variable

$$
\begin{aligned}
T & =\theta(t, \mathbf{q} ; \varepsilon) \\
Q_{k} & =\phi_{k}(t, \mathbf{q} ; \varepsilon)
\end{aligned}
$$

and the infinitesimal generators are

$$
\begin{aligned}
\xi & =\left.\frac{\partial \theta}{\partial \varepsilon}\right|_{\varepsilon=0} \\
\eta_{k} & =\left.\frac{\partial \phi_{k}}{\partial \varepsilon}\right|_{\varepsilon=0}
\end{aligned}
$$

## More than one dependent variable

## Common symmetries

Noether's theorem: Suppose $L(t, \mathbf{q}, \dot{\mathbf{q}})$ is variationally invariant on $\left[t_{0}, t_{1}\right]$ under a transform with infinitesimal generators $\xi$ and $\eta_{k}$. Given

$$
p=\frac{\partial L}{\partial \dot{q}_{k}}, \quad H=\sum_{k=1}^{n} p_{k} \dot{q}_{k}-L
$$

Then

$$
\sum_{k=1}^{n} p_{k} \eta_{k}-H \xi=\text { const }
$$

along any extremal of

$$
F\{\mathbf{q}\}=\int_{t_{0}}^{t_{1}} L(t, \mathbf{q}, \dot{\mathbf{q}}) d t
$$

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## Example: rotations

Invariance in rotations, i.e.

$$
\begin{aligned}
\left(T, Q_{1}, Q_{2}\right) & =\left(t, q_{1} \cos \varepsilon+q_{2} \sin \varepsilon,-q_{1} \sin \varepsilon+q_{2} \cos \varepsilon\right) \\
\left(t, q_{1}, q_{2}\right) & =\left(T, Q_{1} \cos \varepsilon-Q_{2} \sin \varepsilon, Q_{1} \sin \varepsilon+Q_{2} \cos \varepsilon\right)
\end{aligned}
$$

The infinitesimal generators are

$$
\begin{aligned}
\xi & =0 \\
\eta_{1} & =-q_{1} \sin \varepsilon+\left.q_{2} \cos \varepsilon\right|_{\varepsilon=0} \\
\eta_{2} & =-q_{2} \\
\eta_{1} \cos \varepsilon-\left.q_{2} \sin \varepsilon\right|_{\varepsilon=0} & =-q_{1}
\end{aligned}
$$

So, a system with such invariance has

$$
\sum_{i=1}^{2} p_{i} \eta_{i}-H \xi=p_{1} q_{2}-p_{2} q_{1}=\mathrm{const}
$$

So angular momentum in conserved.

Given a system in 3D with Kinetic Energy $T(\dot{\mathbf{q}})=\frac{1}{2} m\left(\dot{q}_{1}^{2}+\dot{q}_{2}^{2}+\dot{q}_{3}^{2}\right)$, and Potential Energy $V(t, \mathbf{q})$.

- invariance of $L$ under time translations corresponds to conservation of Energy
- invariance of $L$ under spatial translations corresponds to conservation of momentum
- invariance of $L$ under rotations corresponds to conservation of angular momentum


## Finding symmetries

Testing for non-trivial symmetries can be tricky.
Useful result is the Rund-Trautman identity:
It leads also to a simple proof of Noether's theorem

## More advanced cases

- Laplace-Runge-Lenz vector in planetary motion corresponds to rotations of 3D sphere in 4D
- symmetries in general relativity
- symmetries in quantum mechanics
- symmetries in fields

