# Variational Methods and Optimal Control Numerical Questions 

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## Some numerical questions and examples:

1. Approximation: In many applications, we want to approximate a piece-wise continuous function $R(x)$ with a smooth curve. For instance:

- it is common to compile histograms of a set of data, but a typical pltted histogram can look artificially "blocky";
- in many signal processing applications we will have a signal (often sampled at discrete time intervals), that contains noise, and one method for removing white noise is to smooth the data.

We can approximate a piecewise constant curve arbitrarily closely, but the cost is the slope of the approximating function may become arbitrarily large. The resulting approximation won't look any smoother than the original.
To obtain a "smoother" approximation, we need to constrain the derivative of the approximation function in some way, but a fixed constraint $\left|y^{\prime}\right| \leq c$ would be rather rigid. Instead let us solve the optimization problem to minimize

$$
F\{y\}=\int_{a}^{b} \gamma^{2}(y-R)^{2}+y^{\prime 2} d x
$$

over the interval $[a, b]$, with $y(x)$ free at each end point, where the parameter $\gamma$ allows a trade off between degree of "smoothness" against a closer fit to the original.
Solution: The Euler-Lagrange equations are

$$
\frac{d}{d x} \frac{\partial f}{\partial y^{\prime}}-\frac{\partial f}{\partial y}=2 y^{\prime \prime}-2 \gamma^{2}(y-R)=0
$$

So to find our approximation we need to find the solution to

$$
y^{\prime \prime}-\gamma^{2} y=-\gamma^{2} R(x)
$$

for piece-wise constant function $R(x)$. The solution to the homogeneous DE $y^{\prime \prime}-\gamma^{2} y=0$ is

$$
y_{h}=A e^{\gamma x}+B e^{-\gamma x} .
$$

Break the interval $[a, b]$ into the segments where $R(x)$ is constant, i.e., take a series of points $a=x_{0}<x_{1}<\cdots<$ $x_{n}=b$ where $R(x)=c_{i}$ for all $x \in\left[x_{i}, x_{i+1}\right)$, then on each segment the DE will have particular solution $y_{i}=c_{i}$ and so on that segment we get

$$
\begin{aligned}
y_{i}(x) & =a_{i} e^{\gamma x}+b_{i} e^{-\gamma x}+c_{i} \\
y_{i}^{\prime}(x) & =a_{i} \gamma e^{\gamma x}-b_{i} \gamma e^{-\gamma x}
\end{aligned}
$$

and the constants $a_{i}$ and $b_{i}$ will be determined by the end-points (which are not known for each segment a priori).
To find the actual approximating curve we have to piece together a series of such segments so that
(a) the end-point values all match so the curve is continuous,
(b) the derivatives match at the end points, so the curve is smooth, and
(c) the natural boundary conditions (where here are equivalent to $y^{\prime}=0$ at the boundary) are also satisfied.

The result will be our smooth interpolation of the piecewise constant curve.
In detail, at each point $x_{i}$ for $i=1,2, \ldots, n-1$ we get a pair of conditions

$$
\begin{aligned}
y_{i-1}\left(x_{i}\right) & =y_{i}\left(x_{i}\right) \\
y_{i-1}^{\prime}\left(x_{i}\right) & =y_{i}^{\prime}\left(x_{i}\right)
\end{aligned}
$$

which is $2(n-1)$ conditions + the two natural boundary conditions where we have $2 n$ unknowns (the $\left(a_{i}, b_{i}\right)$ for $i=0, \ldots, n-1)$. Using the form of the solution $y_{i}$ in our answer we get

$$
\begin{aligned}
a_{i-1} e^{\gamma x_{i}}+b_{i-1} e^{-\gamma x_{i}}-a_{i} e^{\gamma x_{i}}-b_{i} e^{-\gamma x_{i}} & =c_{i}-c_{i-1} \\
\gamma\left[a_{i-1} e^{\gamma x_{i}}-b_{i-1} e^{-\gamma x_{i}}-a_{i} e^{\gamma x_{i}}+b_{i} e^{-\gamma x_{i}}\right] & =0
\end{aligned}
$$

for $i=1,2, \ldots, n-1$, and the natural boundary conditions

$$
\begin{aligned}
a_{0} e^{\gamma a}-b_{0} e^{-\gamma a} & =0 \\
a_{n-1} e^{\gamma b}-b_{n-1} e^{-\gamma b} & =0
\end{aligned}
$$

We have $2 n$ linear equations in $2 n$ unknowns. Solving gives the co-efficients we need to obtain the solution. For an example, see Figure 1, which shows the approximations for three different $\gamma$ values.


Figure 1: Approximations for three different $\gamma$ values.
Remarks: We can obviously generalize this method to deal with approximation of curves that are piecewise linear, or for interpolating a function which we only know at some points.

Moreover, the method above is also a special case of a more general class of smoothing and approximation algorithms. In general, we are seeking to minimise some distance between our approximation and the original function, while also meeting some smoothness criteria. In general we might write this as find the curve $y$ that minimises

$$
F\{y\}=\gamma d_{z}\{y\}+I\{y\}
$$

where $d$ is a distance metric between the original function $z$ and the approximation $y$, and $I(\cdot)$ is a function that measures the irregularity, or lack of smoothness of $y$. In the case above, both of these are integrals:

$$
\begin{aligned}
d_{z}\{y\} & =\int_{a}^{b}(y-z)^{2} d x \\
I\{y\} & =\int_{a}^{b} y^{\prime 2} d x
\end{aligned}
$$

However, we could replace these with alternatives.

- We could measure irregularity by minimizing curvature of the aprpoximating functiun, as measured by

$$
I\{y\}=\int_{a}^{b} y^{\prime \prime 2} d x
$$

or with higher-order derivatives, or different powers, or some combination thereof.

- We can use different powers $p$ in the distance function (corresponding to the $\ell_{p}$-norm distance. We could also measure the distance only on a set of sample points $x_{i}$, e.g., we could take a function

$$
d_{z}\{y\}=\sum_{i}\left(y\left(x_{i}\right)-y_{i}\right)^{2}
$$

- We can replace either of the two optimization criteria with hard constraints, e.g., $\left|y^{\prime}\right| \leq 1$, or $y\left(x_{i}\right)-y_{i}=0$, which can be incorporated into the problem in the usual way.

Particular combinations of these two result in well-known special cases:

- If we require that $y\left(x_{i}\right)=y_{i}$, e.g., we require the curve to match $R(x)$ at the grid points $x_{i}$, and minimize curvature of the approximation should be minimal, then the Euler-Poisson equations are

$$
y^{(3)}=0
$$

subject to the constraints, so the curves will be piece-wise cubic. This leads naturally to interpolating splines.

- Smoothing splines arise if we aim to minimize

$$
F\{y\}=\gamma \sum_{i}\left(y\left(x_{i}\right)-y_{i}\right)^{2}+\int_{a}^{b} y^{\prime \prime 2} d x
$$

As $\gamma$ approaches 1, we approach the interpolative splines above, whereas, when $\gamma$ approaches zero we obtain a linear-regression to the data.

- De Boor extended this to allow for weights on each data point, and higher-order derivatives, e.g.,

$$
F\{y\}=\gamma \sum_{i}\left(\frac{y\left(x_{i}\right)-y_{i}}{\delta_{i}}\right)^{2}+\int_{a}^{b} y^{(m)^{2}} d x
$$

The critical point is that nearly all of these "smoothers" can be derived from the calculus of variations.
The choice of smoothing parameter $\gamma$ is a hard problem in general and is often a little arbitrary, though in specific cases there may be methods to choose it quantitatively, based on known issues such as the level of noise in a dataset.

Matlab code: for performing calculating the co-efficients for an arbitrary piece-wise constant function is included below:

```
function [a,b,c] = approx_function(x_0, x_n, x_i, R_i, gamma)
% file: approx_function.m, (c) Matthew Roughan, Thu Oct 28 2010
created: Thu Oct 28 2010
% author: Matthew Roughan, matthew.roughan@adelaide.edu.au
R is a piece-wise constant function which we wish to approximate
with minimal deviation, and slope, i.e., minimize
    \[F\{ Y \} = \int_{a}^{b} \gamma^2 (y-R)^2 + y`^2 \, dx , \]
input:
        x_0
        x_n = the function R(x) is defined on [x_0, x_n]
        x_i = a nxl vector of the x-values where the function R changes
            x_i(1) = x_0
            but x_n is not included in the vector x_i
        R_i = a nxl vector of the values of R(x) on each sub-interval
        gamma = the tradeoff parameter between fidelity to the function and the
                        allowed slope of the approximation
    outputs:
        the optimal approximation function takes the form
            y_i (x) = a_i e^{\gamma x} + b_i e^{-\gamma x} + c_i
        on each interval [x_i(i), x_i(i+1)]. The outputs of the function are vectors of
        the coefficients a_i and b_i
            a = a nxl vector of the coefficients of a_i in the above approximation
            b = a nxl vector of the coefficients of b_i in the above approximation
            c = a nxl vector of the coefficients of c_i in the above approximation (here c_i = R_i)
n = length(x_i);
% check inputs
sx = size(x_i)
sR = size(R_i)
if (sx(1) == 1& &x(2) > 1)
    x_i = x_i';
elseif (sx(1) > 1 & sx(2) > 1)
    error('sx should be a nx1 vector');
end
if (sR(1) == 1 & sR(2) > 1)
    R_i = R_i';
elseif (sR(1) > 1 & sR(2) > 1)
    error('sR should be a nx1 vector');
end
if (sx(1) ~ = sR(1) | sx(2) ~= sR(2))
    error('x and R should be the same size');
end
if (x_0 ~ = x_i(1))
    error('we should have x_i(1)=x_0');
end
% temporary variables
c = R_i
d = exp(gamma*x_i);
e = exp(-gamma*x_i);
% set up continuity equations at (n-1) joins x_i(2:end)
    a_{i-1} e^{\gamma x_i} + b_{i-1} e^{-\gamma x_i} - a_{i} e^{\gamma x_i} - b_{i} e^{-\gamma x_i}
        matrix A1 are coefficients for the a_i
        matrix B1 are coefficients for the b_i
        y1 are the right-hand side terms of the equations (c_{i} - c_{i-1)
A1 = zeros (n-1,n);
B1 = zeros(n-1,n);
for i=1:n-1
    A1(i,i) = d(i+1);
```

```
    A1(i,i+1) = -d(i+1);
    B1(i,i) = e(i+1);
    B1(i,i+1) = -e(i+1);
end
y1 = [diff(c)];
set up continuity equations for derivatives at (n-1) joins x_i(2:end)
\gamma \left[ a_{i-1} e^{\gamma x_i} - b_{i-1} e^{-\gamma x_i} - a_i e^{\gamma x_i} + b_i e^{-\gamma x_i}
    matrix A2 are coefficients for the a_i
    matrix B2 are coefficients for the b_i
    y2 are the right-hand side terms of the equations (zero in this case)
A2 = zeros(n-1,n);
B2 = zeros (n-1,n);
for i=1:n-1
    A2(i,i) = d(i+1);
    A2(i,i+1) = -d(i+1);
    B2(i,i) = -e(i+1);
    B2(i,i+1) = e(i+1);
end
y2 = zeros(n-1,1);
% complete set of constraints at the joins
A = [[A1, B1]
        [A2, B2]];
y = [y1; y2];
% add in natural boundary conditions at the edges
% natural boundary says y'=0 at x_0 and x_n
A = [[exp (gamma*x_0), zeros(1,n-1), -exp(-gamma*x_0), zeros(1,n-1)];
        A;
        [zeros(1,n-1), exp(gamma*x_n), zeros(1,n-1), -exp(-gamma*x_n)];
        ]
y = [0;y;0]
% solve the equations to find 'coefficients', which consists of
                coefficients = [a; b]
% where a = [a_i], and b=[b_i] and A*coefficients = y
coefficients = A \ y;
a = coefficients(1:n)
b = coefficients(n+1:end)
```


## 2. Catenary:

The shape of a hanging chain of length $L$ was presented as the solution of the problem of minimizing potential energy

$$
W_{p}\{y\}=m g \int_{x_{0}}^{x_{1}} y \sqrt{1+y^{\prime 2}} d x
$$

under the isoperimetric constraint

$$
G\{y\}=\int_{x_{0}}^{x_{1}} \sqrt{1+y^{\prime 2}} d x=L
$$

assuming the (given) heights of the pylons $y_{i}=y\left(x_{i}\right)>0$.
We determined that the solution to this problem took the form

$$
y=c_{1} \cosh \left(\frac{x-c_{2}}{c_{1}}\right)-\lambda
$$

where the constants $\lambda, c_{1}$ and $c_{2}$ are determined by the length $L$ of the chain, and the end conditions, i.e., the heights of the poles $y\left(x_{0}\right)=x_{0}$ and $y\left(x_{1}\right)=x_{1}$.
Determine:
(a) the infimum of possible lengths of chain, and
(b) the maximum length before the chain drags on the ground (at height $y=0$ ).

## Solution:

(a) The minimum length of the chain is not well defined because although the minimum distance (along a straight line) is well defined, a chain cannot take this shape (except in the limit as tension in the chain goes to infinity). Hence we use the term infimum, which can be thought of as the greatest lower bound.
Formally, the infimum of a subset $S$ of some partially ordered set $T$ is the greatest element of $T$ that is less than or equal to all elements of $S$. Here, the set $T$ is the set of all curves $y(x)$ with two continuous derivatives, that satisfy the end-point constraints. The ordering is based on the lengths of the curves, and the subset $S$ is the set of all catenaries. The infimum of the lengths will simply be the length of the straight line between the two end-points, because we know that we can get a catenary arbitraily close to this line. Hence

$$
L_{\mathrm{inf}}=\sqrt{\left(x_{1}-x_{0}\right)^{2}+\left(y_{1}-y_{0}\right)^{2}}
$$

That is, even though we can't have a catenary which is a straight line, it forms the greatest lower bound on the lengths of the catenaries, i.e., the infimum.
(b) We can calculate the maximum length, by noting that the chain will only drag on the ground if its minimum height $y_{\text {min }}<0$, so the maximum chain length will occur when $y_{\min }=0$. Note that

$$
\begin{aligned}
y & =c_{1} \cosh \left(\frac{x-c_{2}}{c_{1}}\right)-\lambda \\
y^{\prime} & =\sinh \left(\frac{x-c_{2}}{c_{1}}\right)
\end{aligned}
$$

which has a zero at $x=c_{2}$. So there are three possible locations for the mimimum - the two edges, or the stationary points $y^{\prime}=0$.

$$
x_{\min }= \begin{cases}x_{0}, & \text { if } c_{2} \leq x_{0}  \tag{1}\\ c_{2}, & \text { if } x_{0} \leq c_{2} \leq x_{1} \\ x_{1}, & \text { if } c_{2} \geq x_{1}\end{cases}
$$

In these cases we get

$$
y_{\min }= \begin{cases}y_{0}, & \text { if } c_{2} \leq x_{0}  \tag{2}\\ c_{1}-\lambda, & \text { if } x_{0} \leq c_{2} \leq x_{1} \\ y_{1}, & \text { if } c_{2} \geq x_{1}\end{cases}
$$

However, in order that we have the maxmium chain length $y_{\text {min }}=0$, so we know that $x_{0} \leq c_{2} \leq x_{1}$, and

$$
c_{1}=\lambda
$$

The form of the catenary is therefore

$$
y=c_{1}\left[\cosh \left(\frac{x-c_{2}}{c_{1}}\right)-1\right]
$$

Note that $\cosh ^{2}(x)-\sinh ^{2}(x)=1$, so

$$
\begin{aligned}
\left.c_{1} \sinh \left(\left(x-c_{2}\right) / c_{1}\right)\right) & =\operatorname{sign}\left(x-c_{2}\right) \sqrt{\left.c_{1}^{2} \cosh ^{2}\left(\left(x-c_{2}\right) / c_{1}\right)\right)-c_{1}^{2}} \\
& =\operatorname{sign}\left(x-c_{2}\right) \sqrt{\left(y_{1}+\lambda\right)^{2}-c_{1}^{2}}
\end{aligned}
$$

Hence

$$
\begin{aligned}
L\{y\} & \left.=c_{1}\left[\sinh \left(\left(x-c_{2}\right) / c_{1}\right)\right)\right]_{x_{0}}^{x_{1}} \\
= & {\left[\operatorname{sign}\left(x-c_{2}\right) \sqrt{(y+\lambda)^{2}-c_{1}^{2}}\right]_{x_{0}}^{x_{1}} } \\
& = \begin{cases}\sqrt{\left(y_{1}+\lambda\right)^{2}-c_{1}^{2}}-\sqrt{\left(y_{0}+\lambda\right)^{2}-c_{1}^{2}}, & \text { if } c_{2}<x_{0}, \\
\sqrt{\left(y_{1}+\lambda\right)^{2}-c_{1}^{2}}, & \text { if } c_{2}=x_{0}, \\
\sqrt{\left(y_{1}+\lambda\right)^{2}-c_{1}^{2}}+\sqrt{\left(y_{0}+\lambda\right)^{2}-c_{1}^{2}}, & \text { if } x_{0}<c_{2}<x_{1}, \\
\sqrt{\left(y_{0}+\lambda\right)^{2}-c_{1}^{2}}, & \text { if } c_{2}=x_{1}, \\
-\sqrt{\left(y_{1}+\lambda\right)^{2}-c_{1}^{2}}+\sqrt{\left(y_{0}+\lambda\right)^{2}-c_{1}^{2}}, & \text { if } c_{2}>x_{1} .\end{cases}
\end{aligned}
$$

but as we know for the maximum length chain that $\lambda=c_{1}$ and $x_{0} \leq c_{2} \leq x_{1}$, we can write

$$
L\{y\}=\sqrt{y_{1}\left(y_{1}+2 c_{1}\right)}+\sqrt{y_{0}\left(y_{0}+2 c_{1}\right)}
$$

We can determine $c_{1}$ and $c_{2}$ now by numerical solution of the end-point equations. We solve by constructing a function which is the square of the deviation from the two end-point constraints: i.e., we define

$$
\begin{aligned}
& g_{1}\left(c_{1}, c_{2}\right)=y_{0}-c_{1} \cosh \left(\frac{x_{0}-c_{2}}{c_{1}}\right)-c_{1} \\
& g_{2}\left(c_{1}, c_{2}\right)=y_{1}-c_{1} \cosh \left(\frac{x_{1}-c_{2}}{c_{1}}\right)-c_{1}
\end{aligned}
$$

We then use Matlab's optimization toolbox function fminsearch to find the minimum of

$$
g\left(c_{1}, c_{2}\right)=g_{1}^{2}+g_{2}^{2}
$$

Then we use these to determine $L_{\text {max }}$. Notice also that as $\cosh (x) \geq 1$ for all $x$, and is equal to 1 only at $x=0$, the point at which the curve touches the ground will be $c_{2}$.

Matlab code: for performing estimating the catenary parameters is included below.

```
function [L_max, L_min, c_1, c_2, lambda] = catenary_max_length(y_0, y_1, x_0, x_1)
%
file: catenary_max_length.m, (c) }2012\mathrm{ Matthew Roughan
author: Matthew Roughan
email: matthew.roughan@adelaide.edu.au
CATENARY_SOLVER: calculate the maximum (and min) length of a hangling chain before it dangles on the
                                    ground (is drawn taught), where it takes catenary shape
                            y = c_1*cosh((x-c_2)/c_1) - c_1
    with fixed length
            L = c_1.*( sinh((x_1-c_2)./c_1) - sinh((x_0-c_2)./c_1) );
    but we need to work out the constants of integration c_1, c_2 and lambda
INPUTS:
    Y_0 = height of the left pylon
    y_1 = height of the right pylon
    x_0 = left pylon position
    x_1 = right pylon position
OUTPUTS:
    L_max = the maximum length of the chain
    L_min = the infimum length of the chain
    c_1,c_2,lambda = parameters of maximal catenary
__min = sqrt( (x_1 - x_0).^2 + (y_1 - y_0).^2 );
create a function which we will minimize to find the solution
        g1 is the left end-point constraint
        g2 is the right end-point constraint
        a = [c_1, c_2], lambda = c_1
g1 = @(a) ( y_0 - a(1)*cosh( (x_0 - a(2))/a(1) ) + a(1) ).^2;
g2 = @(a) ( y_1 - a(1)*cosh( (x_1 - a(2))/a(1) ) + a(1) ).^2;
g = @(a) g1(a) + g2(a);
a_est = [100, (x_0+x_1)/2];
options = optimset('fminsearch');
options = optimset(options, 'MaxFunEvals', 100000, 'MaxIter', 1000);
[a, fval, exitflag, output] = fminsearch(g, a_est, options);
c_1 = a(1); % must be > 0
c_2 = a(2); % must be between x_0 and x_1 for the maximum
lambda = -c_1; % for the maximum
g_val = [g1(a), g2(a)];
% compute the maximum length
L_max = c_1.*( sinh((x_1-c_2)./c_1) - sinh((x_0-c_2)./c_1) );
```


## 3. Catenary (numerical solution of the constants):

The shape of a hanging chain of length $L$ was presented as the solution of the problem of minimizing potential energy

$$
W_{p}\{y\}=m g \int_{x_{0}}^{x_{1}} y \sqrt{1+y^{\prime 2}} d x
$$

under the isoperimetric constraint

$$
G\{y\}=\int_{x_{0}}^{x_{1}} \sqrt{1+y^{\prime 2}} d x=L
$$

We determined that the solution to this problem took the form

$$
y=c_{1} \cosh \left(\frac{x-c_{2}}{c_{1}}\right)-\lambda
$$

where the constants $\lambda, c_{1}$ and $c_{2}$ are determined by the length $L$ of the chain, and the end conditions, i.e., the heights of the poles $y\left(x_{0}\right)=x_{0}$ and $y\left(x_{1}\right)=x_{1}$.
We determined a method to calculate the constants when the problem is symmetric, though this required numerical solution of a non-linear equation. In this more general case, we can make use of identity such as

$$
\begin{aligned}
L\{y\} & =\int_{x_{0}}^{x_{1}} \sqrt{1+y^{\prime 2}} d x \\
& =\int_{x_{0}}^{x_{1}} \cosh \left(\left(x-c_{2}\right) / c_{1}\right) d x \\
& \left.=c_{1}\left[\sinh \left(\left(x-c_{2}\right) / c_{1}\right)\right)\right]_{x_{0}}^{x_{1}}
\end{aligned}
$$

Given a length $L>D$ where $D$ is the distance between the two pylon ends, write code to numerically determine the constants $\left(c_{1}, c_{2}, \lambda\right)$.
Solution: Given a length $L>D$ where $D$ is the distance between the two pylon ends, there will always be a valid catenary, because we know we could hang a chain between these two points (ignoring the possibility that it would drag on the surface).

We solve by constructing a function which is the square of the deviation from the three available constraints: i.e., we define

$$
\begin{aligned}
g_{1}\left(c_{1}, c_{2}, \lambda\right) & =y_{0}-c_{1} \cosh \left(\frac{x_{0}-c_{2}}{c_{1}}\right)-\lambda \\
g_{2}\left(c_{1}, c_{2}, \lambda\right) & =y_{1}-c_{1} \cosh \left(\frac{x_{1}-c_{2}}{c_{1}}\right)-\lambda \\
g_{3}\left(c_{1}, c_{2}, \lambda\right) & \left.=L-c_{1}\left[\sinh \left(\left(x-c_{2}\right) / c_{1}\right)\right)\right]_{x_{0}}^{x_{1}}
\end{aligned}
$$

We then use Matlab's optimization toolbox function fminsearch to find the minimum of

$$
g\left(c_{1}, c_{2}, \lambda\right)=g_{1}^{2}+g_{2}^{2}+g_{3}^{2} .
$$

Matlab code is provided below, as are some example results.


## Matlab code: for performing estimating catenary parameters is included below.

```
function [x, y, c_1, c_2, lambda, Lest, Fest, Lest_check, Fest_check, ...
    f_val, exitflag, output] = catenary_solver_gen(n, y_0, y_1, x_0, x_1, L)
%
    file: catenary_solver_gen.m, (c) 2012 Matthew Roughan
    author: Matthew Roughan
    email: matthew.roughan@adelaide.edu.au
    CATENARY_SOLVER: solves the shape of a hanging chain, which we know will be
                y = c_1*\operatorname{cosh}((x-c_2)/c_1) + lambda
                with fixed length
                    L = c_1.*( sinh((x_1-c_2)./c_1) - sinh((x_0-c_2)./c_1) );
                but we need to work out the constants of integration c_1, c_2 and lambda
    INPUTS:
            n = number of points at which to calculate the curve
            y_0 = height of the left pylon (must be > 0)
            y_1 = height of the right pylon (must be > 0)
            x_0 = left pylon position
            x_1 = right pylon position
            L = length of chain
    OUTPUTS:
            [x,y] = n (x,y) points along the shape of the catenary
            c_1,c_2 = constants of integration
            lambda}=\mathrm{ Lagrange multiplier
            Lest = estimated length, to be used in debugging
            Fest = an estimate of the functional which gives the potential energy of the chain
            Lest_check = a check based on estimated (x,y) positions (only valid for large n)
            Fest_check = a check based on estimated (x,y) positions (only valid for large n)
            [f_val, exitflag, output] = output from the optimization used to find the solution
    f(y_0 < 0)
        error('Y_0 must be >= 0');
end
    f (y_1 < 0)
        error('y_1 must be >= 0');
end
if (x_1 <= x_0)
    error('x_1 should be > x_0');
end
if (nargin == 6) % L defined
    [x, y, c_1, c_2, lambda, Lest, Fest, Lest_check, Fest_check, ...
    f_val, exitflag, output] = catenary_new_b (y_0, y_1, x_0, x_1, L);
elseif (nargin == 5) % compute the natural catenary
    lambda = 0;
    % create a function which we will minimize to find the solution
            g1 is the left end-point constraint
            g2 is the right end-point constraint
    remembering for this case, there may be one, zero, or two solutions
    g1 = @(a) ( y_0 - a(1)*cosh( (x_0 - a(2))/a(1) )).^2;
    g2 =a(a) (y_1 - a(1)*\operatorname{cosh( (x_1 -a(2))/a(1) )).^2;};
    g = @(a) g1(a) + g2(a);
    a_est = [0.1, (x_1+x_0)/2]; % case with smaller c_1 is usually the min
    options = optimset('fminsearch');
    options = optimset(options, 'MaxFunEvals', 100000);
```

```
    [a, fval1, exitflag1, output1] = fminsearch(g, a_est, options);
    c_11 = a(1);
    c_21 = a(2);
    Fest_1 = c_11*(x_1 - x_0)/2 + ...
    c_11^2*(sinh(2*(x_1 - c_21)/c_11)-sinh(2*(x_0 - c_21)/c_11))/4;
    a_est = [100, (x_1+x_0)/2]; % case with larger c_1 is important to check for as well
    options = optimset('fminsearch');
    options = optimset(options, 'MaxFunEvals', 100000);
    [a, fval2, exitflag2, output2] = fminsearch(g, a_est, options);
    c_12 = a(1);
    c_22 = a(2);
    lambda = 0;
    Fest_2 = c_12*(x_1 - x_0)/2 + ...
    c_12^2*(sinh(2*(x_1 - c_22)/c_12)-sinh(2*(x_0 - c_22)/c_12))/4;
    f_val = [fval1, fval2];
if (fval1 > 1.0e-6 & fval2> 1.0e-6)
    % no solutions
    c_1 = NaN;
    c_2 = NaN;
    exitflag=exitflag1;
    output=output1;
elseif (Fest_1 < Fest_2)
    c_1 = c_11;
    c_2 = c_21;
    exitflag=exitflag1;
    output=output1;
else
    c_1 = c_12;
    c_2 = c_22;
    exitflag=exitflag2;
    output=output2;
end
end
% check the length is correct
Lest = c_1.*( sinh((x_1-c_2)./c_1) - sinh((x_0-c_2)./c_1) );
%
% now calculate points on the curve
%
x = x_0:(x_1 - x_0)/n:x_1;
y = c_1*cosh((x-c_2)/c_1) + lambda;
% second check of the length
Lest_check = sum(sqrt(diff(x).^2 + diff(y).^2));
```

```
% calculate the potential a few different ways to compare
```

% calculate the potential a few different ways to compare
Fest_1 = sum(y(1:end-1).*sqrt(diff(x).^2 + diff(y).^2));
Fest_1 = sum(y(1:end-1).*sqrt(diff(x).^2 + diff(y).^2));
Fest_2 = c_1*(x_1 - x_0)/2 + ...
Fest_2 = c_1*(x_1 - x_0)/2 + ...
c_1^2*(sinh(2*(x_1 - c_2)/c_1)-sinh(2*(x_0 - c_2)/c_1))/4 - ...
c_1^2*(sinh(2*(x_1 - c_2)/c_1)-sinh(2*(x_0 - c_2)/c_1))/4 - ...
lambda*Lest;
lambda*Lest;
% Fest_3 = c_1*(x_1 - x_0)/2 + ...
% Fest_3 = c_1*(x_1 - x_0)/2 + ...
c_1*((y_1+lambda)*sinh((x_1 - c_2)/c_1) - (y_0+lambda)*sinh((x_0 - c_2)/c_1))/2 - ...
c_1*((y_1+lambda)*sinh((x_1 - c_2)/c_1) - (y_0+lambda)*sinh((x_0 - c_2)/c_1))/2 - ...
lambda*Lest;
lambda*Lest;
Fest_4 = c_1*(x_1 - x_0)/2 + ...
Fest_4 = c_1*(x_1 - x_0)/2 + ...
c_1*(y_1*sinh((x_1 - c_2)/c_1) - y_0*sinh((x_0 - c_2)/c_1))/2 - ...
c_1*(y_1*sinh((x_1 - c_2)/c_1) - y_0*sinh((x_0 - c_2)/c_1))/2 - ...
lambda*Lest/2;
lambda*Lest/2;
Fest = Fest_2;
Fest = Fest_2;
Fest_check = Fest_1;

```
Fest_check = Fest_1;
```


## Remarks:

and

$$
\begin{aligned}
F\{y\} & =\int_{x_{0}}^{x_{1}}\left(c_{1} \cosh \left(\left(x-c_{2}\right) / c_{1}\right)-\lambda\right) \sqrt{1+\sinh ^{2}\left(\left(x-c_{2}\right) / c_{1}\right)} d x \\
& =\int_{x_{0}}^{x_{1}} c_{1} \cosh ^{2}\left(\left(x-c_{2}\right) / c_{1}\right) d x-\lambda L \\
& =\int_{x_{0}}^{x_{1}} \frac{c_{1}}{2}\left[1+\cosh \left(2\left(x-c_{2}\right) / c_{1}\right)\right] d x-\lambda L \\
& =\frac{c_{1}}{2}\left(x_{1}-x_{0}\right)+\left[\frac{c_{1}^{2}}{4} \sinh \left(2\left(x-c_{2}\right) / c_{1}\right)\right]_{x_{0}}^{x_{1}}-\lambda L \\
& =\frac{c_{1}}{2}\left(x_{1}-x_{0}\right)+\left[\frac{c_{1}^{2}}{2} \sinh \left(\left(x-c_{2}\right) / c_{1}\right) \cosh \left(\left(x-c_{2}\right) / c_{1}\right)\right]_{x_{0}}^{x_{1}}-\lambda L \\
& =\frac{c_{1}}{2}\left(x_{1}-x_{0}\right)+\left[\frac{c_{1}}{2} \sinh \left(\left(x-c_{2}\right) / c_{1}\right)(y+\lambda)\right]_{x_{0}}^{x_{1}}-\lambda L \\
& =\frac{c_{1}}{2}\left(x_{1}-x_{0}\right)+\frac{c_{1}}{2}\left[y(x) \sinh \left(\left(x-c_{2}\right) / c_{1}\right)\right]_{x_{0}}^{x_{1}}-\frac{\lambda L}{2}
\end{aligned}
$$

which were given in lectures.

$$
\sinh (2 x)=2 \sinh (x) \cosh (x), \quad \text { and } \cosh (2 x)=2 \cosh ^{2}(x)-1
$$

If $y_{1}=y_{2}$ we get the symmetric case, which has $c_{2}=\left(x_{1}+x_{2}\right) / 2$, and from above

$$
\begin{aligned}
F\{y\} & =\frac{c_{1}}{2}\left(x_{1}-x_{0}\right)+\frac{c_{1}}{2}\left[y(x) \sinh \left(\left(x-c_{2}\right) / c_{1}\right)\right]_{x_{0}}^{x_{1}}-\frac{\lambda L}{2} \\
& =\frac{c_{1}}{2}\left(x_{1}-x_{0}\right)+\frac{y_{1}-\lambda}{2} L
\end{aligned}
$$

When $\lambda=0$, we get the natural catenary, and so this must give the lowest value for $F\{y\}$ for any $L$, and hence, $\lambda<0$ for all other cases.
The first question we might therefore wish to answer is the rough location of $c_{2}$.
We can just try each of the cases above, and see which one works ...
$y_{\text {min }}=c_{1}-\lambda$ occurs when $x=c_{2}$, and the derivative will be zero - could be in the middle, or the left or right of the interval of interest
Start with $y$ and approximate cosh using first two terms of Taylor series:

$$
\begin{aligned}
\sinh (x) & =x+\frac{x^{3}}{6}+\frac{x^{5}}{120}+\ldots \\
\cosh (x) & =1+\frac{x^{2}}{2}+\frac{x^{4}}{24}+\ldots
\end{aligned}
$$

so that

$$
\begin{aligned}
y+\lambda & =c_{1} \cosh \left(\frac{x-c_{2}}{c_{1}}\right) \\
& =c_{1}+\frac{\left(x-c_{2}\right)^{2}}{2 c_{1}}+\cdots
\end{aligned}
$$

Taking the values at $x_{0}$ and $x_{1}$, and subtracting we get

$$
\begin{aligned}
y_{1}-y_{0} & =\frac{\left(x_{1}-c_{2}\right)^{2}}{2 c_{1}}-\frac{\left(x_{0}-c_{2}\right)^{2}}{2 c_{1}}+\cdots \\
2 c_{1}\left(y_{1}-y_{0}\right) & \simeq\left(x_{1}-c_{2}\right)^{2}-\left(x_{0}-c_{2}\right)^{2} \\
& \simeq\left(x_{1}^{2}-x_{0}^{2}\right)-2\left(x_{1}-x_{0}\right) c_{2} \\
c_{1} & \simeq \frac{\left(x_{1}^{2}-x_{0}^{2}\right)-2\left(x_{1}-x_{0}\right) 2 c_{2}}{2\left(y_{1}-y_{0}\right)} \\
& \simeq \alpha+\beta c_{2}
\end{aligned}
$$

so there is an approximately linear relationship between $c_{1}$ and $c_{2}$. Substitute this into the equation for $L$, and we get a single non-linear equation which we can solve to get an approximate value for $c_{2}$, and hence $c_{1}$ (by the above) and hence $\lambda$ by the edge equations.
This gives us an initial guess as to the values of $c_{1}, c_{2}$ and $\lambda$ which we can refine to get an optimal value.
4. Direct Approximate Solution: One of the classic problems in the CoV is the catenary problem, which we solved analytically early in the course. The problem reduces to minimising potential energy

$$
W_{p}\{y\}=m g \int_{0}^{L} y(s) d s=m g \int_{x_{0}}^{x_{1}} y \sqrt{1+y^{\prime 2}} d x
$$

subject to the end-point conditions (the fixed heights of the pylons), and the length constraint:

$$
G\{y\}=\int_{0}^{L} y(s) d s=\int_{x_{0}}^{x_{1}} \sqrt{1+y^{\prime 2}} d x=L
$$

Ignoring the constant $m g$, we solved this using Lagrange multipliers, instead seeking extremals of

$$
H\{y\}=\int_{x_{0}}^{x_{1}}(y+\lambda) \sqrt{1+y^{\prime 2}} d x
$$

Now we consider the numerical approximation and solution of this problem. We could apply Euler's finite differences, or Ritz, but instead consider a chain of length $L$ made of $n$ pieces of length $\ell$ (where $n \ell=L$ ). Assume each chain link $i$ starts at the point $\left(x_{i}, y_{i}\right)$ and ends at $\left(x_{i+1}, y_{i+1}\right)$, for $i=0,1, \ldots, n-1$, where $y_{i} \geq 0$ and $x_{i} \in\left[x_{a}, x_{b}\right]$. The length constraint on each link means that

$$
\ell^{2}=\left(x_{i+1}-x_{i}\right)^{2}+\left(y_{i+1}-y_{i}\right)^{2}, \quad \forall i=0,1,2, \ldots, n-1
$$

The start and end points of the chain are fixed at $\left(x_{0}, y_{0}\right)=\left(x_{a}, y_{a}\right)$ and $\left(x_{n}, y_{n}\right)=\left(x_{b}, y_{b}\right)$ respectively, and we assume that $D=\sqrt{\left(x_{b}-x_{a}\right)^{2}+\left(y_{b}-y_{a}\right)^{2}}<L$ to allow for a solution.
The objective function that we aim to optimize is to minimize the potential energy of the chain, where the potential of each link will be given by the height of its mid-point, times a constant ( $m g$ ), i.e.,

$$
W=m g \sum_{i=0}^{n-1} \frac{y_{i}+y_{i+1}}{2}
$$

Note that the constant factors are irrelevant, and the sum concertinas so that we can rewrite it

$$
W=\sum_{i=1}^{n-1} y_{i}
$$

where we ommit the values of the end points because these are fixed.
Use a numerical optimization algorithm to find the shape of the curve that solves this problem, and show numerically that it approaches the analytic solution as $n \rightarrow \infty$.

## Solution:

We can solve the problem by including a Lagrange multiplier for each of the constraints, i.e.,

$$
f(\mathbf{x}, \mathbf{y}, \lambda)=\sum_{i=1}^{n-1} y_{i}+\sum_{i=0}^{n-1} \lambda_{i}\left[\ell^{2}-\left(x_{i+1}-x_{i}\right)^{2}-\left(y_{i+1}-y_{i}\right)^{2}\right]
$$

for some Lagrange mutipliers $\lambda_{i}$.
To apply Steepest Descent, starting at some point $\mathbf{v}_{0}$, we need to calculate the gradient. The gradient of this function is derived as follows:

$$
\begin{aligned}
\frac{\partial f}{\partial x_{i}} & =2 \lambda_{i}\left(x_{i+1}-x_{i}\right)-2 \lambda_{i-1}\left(x_{i}-x_{i-1}\right) \\
\frac{\partial f}{\partial y_{i}} & =1+2 \lambda_{i}\left(y_{i+1}-y_{i}\right)-2 \lambda_{i-1}\left(y_{i}-y_{i-1}\right) \\
\frac{\partial f}{\partial \lambda_{i}} & =\left[\ell^{2}-\left(x_{i+1}-x_{i}\right)^{2}-\left(y_{i+1}-y_{i}\right)^{2}\right]
\end{aligned}
$$

except for $i=0$ and $i=n$, where the end points are fixed so $\partial f / \partial x_{i}=\partial f / \partial y_{i}=0$. The Hessian matrix is hence almost tridiagonal, though we don't need this for the Steepest Descent Algorithm.
Take the case with $n$ even, and then the starting point can have the chain links in a zig-zag. Construct it as follows. Draw the line interpolating between $\left(x_{a}, y_{a}\right)$ and $\left(x_{b}, y_{b}\right)$, i.e.,

$$
y=\frac{y_{b}-y_{a}}{x_{b}-x_{a}}\left(x-x_{a}\right)+y_{a}
$$

which has length

$$
D=\sqrt{\left(x_{b}-x_{a}\right)^{2}+\left(y_{b}-y_{a}\right)^{2}}
$$

Now distribute a set of points $x_{i}^{\prime}$ evenly along a line of length

$$
x_{i}^{\prime}=i D / n
$$

and alternate the $y_{i}^{\prime}$ positions to satisfy the length constraint on each segment:

$$
y_{i}^{\prime}= \begin{cases}y_{a}, & \text { if } i=2 k \\ y_{a}+d & \text { if } i=2 k+1,\end{cases}
$$

where $d$ is chosen to enforce the length constraint on each segment, i.e.,

$$
d=\sqrt{\ell^{2}-(D / n)^{2}}
$$

The points $\left(x_{i}^{\prime} \cdot y_{i}^{\prime}\right)$ are then rotated by $\theta$ about $\left(x_{a}, y_{a}\right)$ to meet the other end-point constraint $\left(x_{b}, y_{b}\right)$, where the rotation takes the form

$$
\binom{x_{i}}{y_{i}}=\left(\begin{array}{cc}
\cos (\theta) & \sin (\theta) \\
-\sin (\theta) & \cos (\theta)
\end{array}\right)\binom{x_{i}^{\prime}-x_{a}}{y_{i}^{\prime}-y_{a}}+\binom{x_{a}}{y_{a}}
$$

where

$$
\begin{aligned}
\cos (\theta) & =\frac{x_{b}-x_{a}}{D} \\
\sin (\theta) & =\frac{y_{b}-y_{a}}{D}
\end{aligned}
$$

Matlab code: for performing the steepest descent is provided below:


Figure 3: .

## 5. Rayleigh-Ritz:

Consider the following problem - find the extremal curves of the following functional:

$$
F\{y\}=\int_{0}^{1}{y^{\prime}}^{2}-2 g(x) y d x
$$

for twice differentiable function $g(\cdot)$, and $y(0)=y(1)=0$.
(a) Solve this problem using the Euler-Lagrange equations in general, and in detail for the special cases
i. $g(x)=\sin (2 \pi x)$.
ii. $g(x)=x$.
(b) Now use the Rayleigh-Ritz method to solve the same problems using the approximation

$$
\tilde{y}_{N}=b_{0}+\sum_{n=1}^{N} a_{n} \sin (2 \pi n x)+b_{n} \cos (2 \pi n x)
$$

## Solution:

(a) The Euler-Lagrange equations are

$$
\frac{d}{d x} \frac{\partial f}{\partial y^{\prime}}-\frac{\partial f}{\partial y}=\frac{d}{d x} 2 y^{\prime}+2 g(x)=2 y^{\prime \prime}+2 g(x)=0
$$

The solution is just to integrate twice, resulting in

$$
\begin{aligned}
y^{\prime} & =-\int g(u) d u+c_{1} \\
y & =-\iint g(u) d u d v+c_{1} x+c_{2}
\end{aligned}
$$

where $c_{1}$ and $c_{2}$ are chosen to satisfy $y(0)=y(1)=0$.
i. So for the particular case $g(x)=\sin (2 \pi x)$, where the double integral will be $-\sin (2 \pi x) / 4 \pi^{2}$, which is zero at the end-points, the extremal takes the form

$$
y(x)=\frac{1}{4 \pi^{2}} \sin (2 \pi x)
$$

ii. For the case $g(x)=x$, the double integral will be $x^{3} / 6$, so the extremals are

$$
y(x)=-\frac{x^{3}}{6}+c_{1} x+c_{2}
$$

Using the end-point conditions $y(0)=y(1)=0$ we get

$$
y(x)=-\frac{x^{3}}{6}+\frac{x}{6}=\frac{x\left(1-x^{2}\right)}{6}
$$

(b) Given the approximation, we get

$$
\begin{aligned}
\tilde{y}_{N} & =b_{0}+\sum_{n=1}^{N} a_{n} \sin (2 \pi n x)+b_{n} \cos (2 \pi n x), \\
\tilde{y}_{N}^{\prime} & =2 \pi \sum_{n=1}^{N} n a_{n} \cos (2 \pi n x)-n b_{n} \sin (2 \pi n x), \\
\tilde{y}_{N}^{\prime} & =4 \pi^{2} \sum_{n=1}^{N} \sum_{m=1}^{N} n m\left[a_{n} a_{m} \cos (2 \pi n x) \cos (2 \pi m x)-a_{n} b_{m} \cos (2 \pi n x) \sin (2 \pi m x)+b_{n} b_{m} \sin (2 \pi n x) \sin (2 \pi m x)\right] .
\end{aligned}
$$

When we integrate, we can use

$$
\begin{aligned}
\int_{0}^{1} \cos (2 \pi n x) \cos (2 \pi m x) d x & =\frac{1}{2} \delta_{m n} \\
\int_{0}^{1} \sin (2 \pi n x) \sin (2 \pi m x) d x & =\frac{1}{2} \delta_{m n} \\
\int_{0}^{1} \sin (2 \pi n x) \cos (2 \pi m x) d x & =0
\end{aligned}
$$

So the "cross-terms" in the above integrals vanish, and we are left with

$$
\begin{aligned}
F\{y\} & =\int_{0}^{1} y^{\prime 2}+2 g(x) y d x \\
F(\mathbf{a}, \mathbf{b}) & =2 \pi^{2} \sum_{n=1}^{N} n^{2}\left(a_{n}^{2}+b_{n}^{2}\right)-2 \sum_{n=1}^{N} a_{n} \int_{0}^{1} g(x) \sin (2 \pi n x) d x-2 \sum_{n=0}^{N} b_{n} \int_{0}^{1} g(x) \cos (2 \pi n x) d x \\
& =2 \pi^{2} \sum_{n=1}^{N} n^{2}\left(a_{n}^{2}+b_{n}^{2}\right)-\sum_{n=1}^{N} a_{n} A_{n}-\sum_{n=0}^{N} b_{n} B_{n},
\end{aligned}
$$

where $a_{n}=2 \int g(x) \sin (2 \pi n x) d x$ and $B_{n}=2 \int g(x) \cos (2 \pi n x) d x$ are the Fourier series coeeficients for $g(x)$ (assuming we repeat the function on $[0,1]$ periodically). We can now take the derivatives, and for $n=1,2, \ldots, N$ we get

$$
\begin{aligned}
\frac{d}{d a_{n}} F(\mathbf{a}, \mathbf{b}) & =4 \pi^{2} n^{2} a_{n}-A_{n} \\
& =0 \\
\frac{d}{d b_{n}} F(\mathbf{a}, \mathbf{b}) & =4 \pi^{2} n^{2} b_{n}-B_{n} \\
& =0
\end{aligned}
$$

or

$$
\begin{aligned}
a_{n} & =\frac{A_{n}}{4 \pi^{2} n^{2}}, \\
b_{n} & =\frac{B_{n}}{4 \pi^{2} n^{2}} .
\end{aligned}
$$

Note that the $b_{n}$ terms don't depend on $n$, so we simply expand until we have enough of these terms to converge. We don't get a condition for $b_{0}$, but we have the end-points conditions to consider, and given $\cos (0)=\cos (2 \pi)=$ 1 , and we require $y(0)=y(1)=0$, we choose

$$
b_{0}=-\sum_{n=1}^{N} b_{n}
$$

In the specific cases of interest:
i. When $g(x)=\sin (2 \pi x)$, the Fourier coefficients are

$$
A_{n}=\left\{\begin{array}{ll}
1, & \text { for } n=1, \\
0, & \text { otherwise },
\end{array} \quad \text { and } \quad B_{n}=0\right.
$$

so for all $n$ we get the exact solution:

$$
\tilde{y}_{n}(x)=\frac{1}{4 \pi^{2}} \sin (2 \pi x)
$$

ii. For the case $g(x)=x$, we consider the Fourier coefficients for the "saw-tooth" function, i.e.,

$$
\begin{aligned}
A_{n} & =2 \int_{0}^{1} x \sin (2 \pi n x) d x=-\frac{1}{\pi n}, \\
B_{n} & =2 \int_{0}^{1} x \cos (2 \pi n x) d x=-\frac{1}{2 \pi n^{2}} \text { actually } 0, \\
\tilde{y}_{n}(x)=b_{0} & +\sum_{n=1}^{n} \frac{1}{4 \pi^{2} n^{3}} \sin (2 \pi n x) .+\sum_{n=1}^{n} \frac{1}{8 \pi^{2} n^{3}} \cos (2 \pi n x) .
\end{aligned}
$$

The following figure shows this approximation for $N=1,2,3, \ldots$ versus the actual extremal curve: If we just took a Fourier expansion of $x\left(1-x^{2}\right)$ we get

$$
\begin{aligned}
\int_{0}^{1} x\left(1-x^{2}\right) d x & =1 / 4 \\
a_{n}=2 \int_{0}^{1} x\left(1-x^{2}\right) \sin (2 \pi n x) d x & =-\frac{3}{2 \pi^{3} n^{3}} \\
b_{n}=2 \int_{0}^{1} x\left(1-x^{2}\right) \cos (2 \pi n x) d x & =-\frac{3}{2 \pi^{2} n^{2}}
\end{aligned}
$$

## 6. Another Problem (not finished yet):

Consider the following problem - find the extremal curves of the following functional:

$$
F\{y\}=\int_{0}^{1} y^{\prime 2}-y^{2} d x
$$

for $y(0)=1$ and $y(1)=e$.
(a) Solve this problem using the Euler-Lagrange equations.
(b) Now use the Rayleigh-Ritz method to solve the problem using the trial functions

$$
\tilde{y}=\sum_{i=0}^{\infty} a_{i} x^{i}
$$

Derive a set of conditions on the $a_{i}$ and use your solution to the previous part to solve these conditions.

## Solution:

(a) The LHS of the Euler-Lagrange equation is

$$
\frac{d}{d x} \frac{\partial f}{\partial y^{\prime}}-f=\frac{d}{d x}\left(2 y^{\prime}\right)-2 y
$$

So the DE is

$$
y^{\prime \prime}-y=0
$$

This is a linear, homogeneous DE, with characteristics equation $m^{2}-1=0$, so $m= \pm 1$, and therefore the solution is in the form

$$
y(x)=A e^{t}+B e^{-t} .
$$

From the end-point conditions $y(0)=1$ and $y(1)=$ we get the constraints

$$
\begin{aligned}
A+B & =1 \\
A e-B / e & =e
\end{aligned}
$$

which has solution $A=1, B=0$, and therefore

$$
y(x)=e^{t}
$$

(b) Use the an approximation function

$$
\begin{aligned}
& y_{n}(x)=[1+(e-1) x]+\sum_{i=1}^{n} a_{i} x^{i}(1-x) . \\
& y_{n}(x)=[1+(e-1) x]+\sum_{i=1}^{n} a_{i} \sin (2 \pi i x) .
\end{aligned}
$$

i. Given $\tilde{y}=\sum_{i=0}^{\infty} a_{i} x^{i}$ we get

$$
\begin{aligned}
y^{2} & =\sum_{i, j=0}^{\infty} a_{i} a_{j} x^{i+j}, \\
y^{\prime 2} & =\sum_{i, j=1}^{\infty} i j a_{i} a_{j} x^{i+j-2} \\
& =\sum_{i, j=0}^{\infty}(i+1)(j+1) a_{i+1} a_{j+1} x^{i+j},
\end{aligned}
$$

so the integral can be written

$$
\begin{aligned}
F\{y\} & =\int_{0}^{1} \tilde{y}^{\prime 2}-\tilde{y}^{2} d x, \\
& =\int_{0}^{1} \sum_{i, j=0}^{\infty}(i+1)(j+1) a_{i+1} a_{j+1} x^{i+j}-\sum_{i, j=0}^{\infty} a_{i} a_{j} x^{i+j} d x, \\
& =\sum_{i, j=0}^{\infty}\left[(i+1)(j+1) a_{i+1} a_{j+1}-a_{i} a_{j}\right] \int_{0}^{1} x^{i+j} d x, \\
& =\sum_{i, j=0}^{\infty} \frac{(i+1)(j+1) a_{i+1} a_{j+1}-a_{i} a_{j}}{i+j+1} \\
& =\sum_{i, j: i \neq j}^{\infty} \frac{(i+1)(j+1) a_{i+1} a_{j+1}-a_{i} a_{j}}{i+j+1}+\sum_{i=0}^{\infty} \frac{(i+1)^{2} a_{i+1}^{2}-a_{i}^{2}}{2 i+1} .
\end{aligned}
$$

Now to find the minimum, we set the partial derivatives to zero and we therefore get (for $i \geq 1$ )

$$
\begin{align*}
\frac{\partial F(\mathbf{a})}{\partial a_{i}} & =\sum_{j: j \neq i}^{\infty} \frac{i(j+1) a_{j+1}}{i+j}+\sum_{j: j \neq i}^{\infty} \frac{-a_{j}}{i+j+1}+\frac{i^{2} a_{i}}{i}+\frac{-a_{i}}{i+1}  \tag{3}\\
& =\sum_{j=0}^{\infty} \frac{i(j+1) a_{j+1}}{i+j}+\sum_{j=0}^{\infty} \frac{-a_{j}}{i+j+1}+\frac{i^{2} a_{i}}{2(i-1)}-\frac{a_{i}}{2(i+1)}  \tag{4}\\
& =\sum_{j=0}^{\infty} \frac{i(j+1) a_{j+1}}{i+j}+\sum_{j=0}^{\infty} \frac{-a_{j+1}}{i+j+2}-\frac{a_{0}}{i+2}+\frac{i^{2} a_{i}}{2(i-1)}+\frac{-a_{i}}{2 i}=0 \tag{5}
\end{align*}
$$

We also know from the end-point conditions that

$$
a_{0}=1 \quad \text { and } \quad \sum_{i} a_{i}=e .
$$

This represents a set of relationships between the $a_{i}$, which might be quite hard to solve, however, we already know that there is a solution in this form:

$$
y=e^{t}=\sum_{i=0}^{\infty} \frac{x^{i}}{i!},
$$

so if we take $a_{i}=1 / i$ ! we should get a solution.
First note that with the above choice $a_{0}=1$, and

$$
\sum_{i=0}^{\infty} a_{i}=\sum_{i=0}^{\infty} 1 / i!=e
$$

Now substitute these $a_{i}$ into the relation (5) and we get
ii. Do you get the same solution?
iii. Why did we choose the particular values of $a_{0}$ and $a_{1}$ as the first terms of the sequence?

## 7. More on Rayleigh-Ritz:

In R-R we use an approximation, typically of the form

$$
y_{n}(x)=\phi_{0}(x)+\sum_{i=1}^{n} c_{i} \phi_{i}(x)
$$

to approximate the function

$$
F\{y\}=\int_{x_{0}}^{x_{1}} f\left(x, y, y^{\prime}\right) d x
$$

by

$$
F\left(c_{1}, c_{2}, \ldots, c_{n}\right)=\int_{x_{0}}^{x_{1}} f\left(x, y_{n}, y_{n}^{\prime}\right) d x
$$

where $\phi_{0}$ satisfies the boundary conditions on $y$ and $\phi_{i}\left(x_{i}\right)=0$.
Show ...

## Solution:

The Leibniz integral rule says that for $f$ and $\partial f / \partial x$ continuous over the region $\left[x_{0}, x_{1}\right] \times\left[y_{0}, y_{1}\right]$ we can differentiate an integral in the following form:

$$
\begin{aligned}
& \frac{d}{d x} \int_{y_{0}}^{y_{1}} f(x, y) d y=\int_{y_{0}}^{y_{1}} \frac{\partial}{\partial x} f(x, y) d y . \\
& F\left(c_{1}, c_{2}, \ldots, c_{n}\right)=\int_{x_{0}}^{x_{1}} f\left(x, y_{n}, y_{n}^{\prime}\right) d x \\
&=\int_{x_{0}}^{x_{1}} f\left(x, \phi_{0}+\sum_{i=1}^{n} c_{i} \phi_{i}, \phi_{0}^{\prime}+\sum_{i=1}^{n} c_{i} \phi_{i}^{\prime}\right) d x \\
& \frac{\partial}{\partial c_{i}} F\left(c_{1}, c_{2}, \ldots, c_{n}\right)=\int_{x_{0}}^{x_{1}} \frac{\partial}{\partial c_{i}} f\left(x, \phi_{0}+\sum_{i=1}^{n} c_{i} \phi_{i}, \phi_{0}^{\prime}+\sum_{i=1}^{n} c_{i} \phi_{i}^{\prime}\right) d x \\
&=\int_{x_{0}}^{x_{1}} \frac{\partial f}{\partial y} \phi_{i}+\frac{\partial f}{\partial y^{\prime}} \phi_{i}^{\prime} d x \\
&=\int_{x_{0}}^{x_{1}}\left[\frac{\partial f}{\partial y}-\frac{d}{d x} \frac{\partial f}{\partial y^{\prime}}\right] \phi_{i} d x, \\
&=0, \forall i=1,2, \ldots, n .
\end{aligned}
$$

using the chain rule and integration by parts (noting that $\phi_{i}\left(x_{i}\right)=0$ for $i=1,2, \ldots, n$ ). Obviously the above would be satisfied if the Euler-Lagrange equations were satisfied. However, the underlying assumption of numerical methods is that the Euler-Lagrange equations are difficult.
When we substitute $y_{n}$ into the above formula we get

$$
\begin{aligned}
\frac{\partial}{\partial c_{i}} F\left(c_{1}, c_{2}, \ldots, c_{n}\right) & =\int_{x_{0}}^{x_{1}}\left[\frac{\partial f}{\partial y}-\frac{d}{d x} \frac{\partial f}{\partial y^{\prime}}\right] \phi_{i} d x \\
& =\int_{x_{0}}^{x_{1}} g\left(x, \phi_{1}, \ldots, \phi_{n}, \phi_{1}^{\prime}, \ldots, \phi_{n}^{\prime}\right) \phi_{i}(x) d x \\
& =0, \forall i=1,2, \ldots, n
\end{aligned}
$$

and if this can be calculated, we can then determine values of $c_{i}$ to approximate the extremal. However, Galerkin noted that an alternative
8. Rayleigh-Ritz: Find approximate solutions to the following DE

$$
\frac{d^{2} u}{d t^{2}}=36 t^{2}+12 t-4
$$

subject to the boundary conditions $u(1)=u(-1)=0$, using Rayleigh-Ritz by constructing an appropriate functional, and approximating $u$ using

$$
u_{n}=\sum_{i=1}^{n} c_{i} t^{i-1}\left(t^{2}-1\right)
$$

for $n=1,2$ and 3 , and comment on the accuracy of these solutions.
Solution: The true solution to this is

$$
u(t)=3 t^{4}+2 t^{3}-2 t^{2}-2 t-1
$$

We can construct a functional for this problem

$$
F\{u\}=\int_{-1}^{1} \frac{1}{2}\left(\frac{d u}{d t}\right)^{2}+\left(36 t^{2}+12 t-4\right) u d x
$$

The Euler-Lagrange DE is the equation to be solved.
Given the approximation:

$$
\begin{aligned}
& u_{n}=\sum_{i=1}^{n} c_{i} t^{i-1}\left(t^{2}-1\right) \\
& u_{n}^{\prime}=\sum_{i=2}^{n}(i-1) c_{i} t^{i-2}\left(t^{2}-1\right)+\sum_{i=1}^{n} 2 c_{i} t^{i}
\end{aligned}
$$

Use Matlab's symbolic manipulation toolkit, we construct the functional for these cases, and then differentiate, to obtain a set of equations for the $c_{i}$.
CODE
The results are:

- $n=1$ :
- $n=2$ :
- $n=3$ :

